

Atoms with two electrons

The Schrödinger equation for a 3-body system

- The He-atom (or an ion with two electrons)
 - Two electrons + a nucleus with charge $+Ze$
 - \Rightarrow a 3-body problem
- Exact, analytic solutions are not possible
- We will need approximation methods
 - Perturbation theory

The Schrödinger equation

- In centre-of-mass coordinates:

$$\left(-\frac{\hbar^2}{2\mu} \nabla_{r_1}^2 - \frac{\hbar^2}{2\mu} \nabla_{r_2}^2 - \frac{\hbar^2}{M} \nabla_{r_1} \cdot \nabla_{r_2} - \frac{Ze^2}{4\pi\epsilon_0 r_1} - \frac{Ze^2}{4\pi\epsilon_0 r_2} + \frac{e^2}{4\pi\epsilon_0 r_{12}} \right) \psi(\vec{r}_1, \vec{r}_2) = E\psi(\vec{r}_1, \vec{r}_2)$$

here: $r_{12} = |\vec{r}_1 - \vec{r}_2|$

Simplifications: $M = \infty \Rightarrow \mu = m_e$

- We introduce Atomic units

Atomic units

- Atomic units (a.u) are used to simplify calculations
 - Most constants disappear from Hamiltonians and the Schrödinger
- Starting point : the following natural constants are set to one:

$$e = m_e = \hbar = \frac{1}{4\pi\epsilon_0} = 1$$

- Be careful with quantitative calculations

Quantity	atomic unit	value in SI units
Charge	e	$1.602\,176\,565 \times 10^{-19}$ C
Mass	m_e	$9.109\,382\,91 \times 10^{-31}$ kg
Angular momentum	\hbar	$1.054\,571\,726 \times 10^{-34}$ J s
Length	a_0	$0.529\,177\,210\,92 \times 10^{-10}$ m
Energy	E_h	$4.359\,744\,34 \times 10^{-18}$ J
Time	\hbar / E_h	$2.418\,884\,326\,502 \times 10^{-17}$ s
Force	E_h / a_0	$8.238\,722\,78 \times 10^{-8}$ N
Velocity	$c \alpha$	$2.187\,691\,263\,79 \times 10^6$ m s ⁻¹
Momentum	\hbar / a_0	$1.992\,851\,740 \times 10^{-24}$ kg m s ⁻¹
Charge density	e / a_0^3	$1.081\,202\,338 \times 10^{12}$ C m ⁻³
Electric potential	E_h / e	27.211 385 05 V
Electric field	$E_h / (e a_0)$	$5.142\,206\,52 \times 10^{11}$ V m ⁻¹
Electric dipole moment	$e a_0$	$8.478\,353\,26 \times 10^{-30}$ C m
Magnetic flux density	$\hbar / (e a_0^2)$	$2.350\,517\,464 \times 10^5$ T

- The 2-electron system Hamiltonian in atomic units:

$$\left(-\frac{\nabla_{r_1}^2}{2} - \frac{\nabla_{r_2}^2}{2} - \frac{Z}{r_1} - \frac{Z}{r_2} + \frac{1}{r_{12}} \right) \psi(\vec{r}_1, \vec{r}_2) = E\psi(\vec{r}_1, \vec{r}_2)$$

- Consequences of the term : $\propto \frac{1}{r_{12}}$
 - $\psi(\vec{r}_1, \vec{r}_2)$ cannot be factorised
 - The exact solutions must be entangled states

Symmetry

- With two electrons, symmetry becomes important
 - Spin will matter, due to symmetry
- Total wave function: product of spatial and spin parts

$$\Psi(q_1, q_2) = \psi(\vec{r}_1, \vec{r}_2) \chi(\vec{s}_1, \vec{s}_2)$$

- The Hamiltonian does not depend on spin
 - the wave function can be factorised

The Pauli principle

- The total wave function for two identical fermions is antisymmetric with respect to exchange of the particles
- Two identical fermions cannot occupy the same quantum state simultaneously
- For the product function $\Psi(q_1, q_2)$, we have two options:
 - $\psi(\vec{r}_1, \vec{r}_2)$ symmetric and $\chi(\vec{s}_1, \vec{s}_2)$ anti-symmetric
 - $\psi(\vec{r}_1, \vec{r}_2)$ anti-symmetric and $\chi(\vec{s}_1, \vec{s}_2)$ symmetric

Exchange symmetry

- The exchange operator: P_{12}

$$P_{12} \Psi(q_1, q_2) = \Psi(q_2, q_1)$$

- A permutation of spatial coordinates:
 - if $P_{12} \Psi(q_1, q_2) = \lambda \Psi(q_1, q_2)$
 - $P_{12}^2 \Psi(q_1, q_2) = \lambda^2 \Psi(q_1, q_2) = \Psi(q_1, q_2)$
 - $\Rightarrow \lambda = \pm 1$

Spin wave functions

- For each of the two spin functions, there are only two options
 - Spin-up or Spin-down

- We define kets in the two spin-spaces:

$$\left\{ \begin{array}{l} |+\rangle_1 = \begin{pmatrix} 1 \\ 0 \end{pmatrix}_1 \\ |-\rangle_1 = \begin{pmatrix} 0 \\ 1 \end{pmatrix}_1 \end{array} \right. \quad \left\{ \begin{array}{l} |+\rangle_2 = \begin{pmatrix} 1 \\ 0 \end{pmatrix}_2 \\ |-\rangle_2 = \begin{pmatrix} 0 \\ 1 \end{pmatrix}_2 \end{array} \right.$$

- Compound spin function ;

- four possibilities:

$$\chi_1(\vec{s}_1, \vec{s}_2) : |+\rangle_1 \otimes |+\rangle_2 = |++\rangle$$

$$\chi_2(\vec{s}_1, \vec{s}_2) : |+\rangle_1 \otimes |-\rangle_2 = |+-\rangle$$

$$\chi_3(\vec{s}_1, \vec{s}_2) : |-\rangle_1 \otimes |+\rangle_2 = |-+\rangle$$

$$\chi_4(\vec{s}_1, \vec{s}_2) : |-\rangle_1 \otimes |-\rangle_2 = |--\rangle$$

- (assume that the spatial functions are different, so the Pauli principle does not forbid χ_1 and χ_4)

- There are 2 problems with this basis:

- Problem 1 :

- χ_1 and χ_4 are exchange symmetric , BUT
- χ_2 and χ_3 are neither symmetric, nor anti-symmetric

- We need a description for the “total spin”
 - In absence of spin-spin interaction: $[\vec{S}_1, \vec{S}_2] = 0$
 - \Rightarrow Logical choice : $\vec{S} \equiv \vec{S}_1 + \vec{S}_2$
 - $\Rightarrow \begin{cases} S_z = S_{1z} + S_{2z} \\ S^2 = S_1^2 + S_2^2 + 2\vec{S}_1 \cdot \vec{S}_2 \end{cases}$
 - \Rightarrow Quantum numbers S and M_S
 - The action of S^2 and S_z on $\chi_1, \chi_2, \chi_3, \chi_4$ can be calculated (using the Pauli spin matrices)

- Problem 2 :
 - $S_z |++\rangle = |++\rangle$
 - $S^2 |++\rangle = 2 |++\rangle$
 - $S_z |+-\rangle = 0$
 - $S^2 |+-\rangle = |+-\rangle + |-+\rangle$
 - $S_z |-+\rangle = 0$
 - $S^2 |-+\rangle = |+-\rangle + |-+\rangle$
 - $S_z |--\rangle = -|--\rangle$
 - $S^2 |--\rangle = 2 |--\rangle$
 - χ_2 and χ_3 are not eigenfunctions to S^2
 - \Rightarrow To have a diagonal basis, where all basis functions are either symmetric or anti-symmetric at exchange, we need to replace χ_2 and χ_3

- New functions:
 - $|S\rangle \propto |+-\rangle + |-+\rangle$
 - $|A\rangle \propto |+-\rangle - |-+\rangle$

- \Rightarrow A basis of four functions: $|S M_S\rangle$
- 3 symmetric functions (a triplet) :
 - $\begin{cases} |1, 1\rangle & = |++\rangle \\ |1, 0\rangle & = |S\rangle = \frac{1}{\sqrt{2}}(|+-\rangle + |-+\rangle) \\ |1, -1\rangle & = |--\rangle \end{cases}$
- 1 anti-symmetric function (a singlet) :
 - $|0, 0\rangle = |A\rangle = \frac{1}{\sqrt{2}}(|+-\rangle - |-+\rangle)$

The ground state of He

Perturbation Theory

- Assume that the interaction term can be treated as a perturbation:

$$H = H_0 + H'$$

$$H_0 = -\frac{\nabla_{r_1}^2}{2} - \frac{\nabla_{r_2}^2}{2} - \frac{Z}{r_1} - \frac{Z}{r_2}$$

$$H' = \frac{1}{r_{12}}$$

- The zero-order solution can be factorized

$$H_0 \psi^{(0)}(\vec{r}_1, \vec{r}_2) = E_0 \psi^{(0)}(\vec{r}_1, \vec{r}_2)$$

$$\psi^{(0)}(\vec{r}_1, \vec{r}_2) = \psi_1^{(0)}(\vec{r}_1) \psi_2^{(0)}(\vec{r}_2)$$

$$E_0 = E_1 + E_2$$

- The zero order ground state will be both electron is hydrogenic 1s-orbitals, with $Z = 2$

$$\psi^{(0)} = \psi_{1s} \psi_{1s} = (R_{1s} Y_{00}) (R_{1s} Y_{00}) = \psi_{1s^2}$$

Identical electrons - the Pauli principle

- We have
 - $n_1 = n_2 = 1$, $l_1 = l_2 = 0$, $m_{l1} = m_{l2} = 0$
 - The compound spatial wave function HAS to be symmetric

- To avoid violation the Pauli principle, the two spins **HAVE** to be opposite

$$\begin{aligned}\Rightarrow \Psi^{(0)} &= \psi_{1s^2} \chi_{0,0} \\ |(0)\rangle &= |1s^2\rangle \otimes |0,0\rangle = |1s^2, 00\rangle \\ &\Rightarrow 1s^2\ ^1S \Rightarrow 1s^2\ ^1S\end{aligned}$$

The energy of the ground state

- We define this as the ionization energy E_{ion}
- The zero-order energy (without the perturbation) :
 - $E_0 = E_1 + E_2 = 2E_{1s}(Z = 2) = 2(-Z^2hcR_\infty)$

$$\approx 2(-54.4 \text{ eV}) \approx -109 \text{ eV}$$

- The perturbation :

$$\begin{aligned}\Delta E &= \langle \psi_{1s^2} | H' | \psi_{1s^2} \rangle \\ H' &= \frac{1}{r_{12}} \\ \psi_{1s^2} &= \left[\left(\frac{Z}{a_m u} \right)^{3/2} 2e^{-\rho} \right]^2 \left[\sqrt{\frac{1}{4\pi}} \right]^2\end{aligned}$$

$$\Rightarrow \Delta E \approx 34 \text{ eV}$$

$$\Rightarrow E(1s^2) \approx -109 \text{ eV} + 34 \text{ eV} = -75 \text{ eV}$$

- This means that 75 eV is the energy needed to remove **BOTH** electrons from the nucleus
 - Suppose one electron has already been removed; how much energy is needed to remove the other one?
 - \Rightarrow the ionization energy of He^+

- $E_{\text{ion}}(\text{He}^+) = E_{1s}(Z = 2) \approx -54.4 \text{ eV}$
- The ionization energy of He :
$$E_{\text{ion}}(\text{He}) = E(1s^2) - E_{\text{ion}}(\text{He}^+) \approx -21 \text{ eV}$$
- Experimental value of the helium ionization energy :
-24.6 eV
 - The order of magnitude is right, but
 - The energy contribution from the electron-electron interaction is too great to be treated as a perturbation

Excited states of He

- One of the electrons is in the 1s-orbital
- The other in an nl -orbital ($n \neq 1$)

$$1s nl$$

Exchange degeneracy

- We have two states with the same energy:

$$\psi_{1s}(\vec{r}_1) \psi_{nl}(\vec{r}_2)$$

and

$$\psi_{nl}(\vec{r}_1) \psi_{1s}(\vec{r}_2)$$

- This is the “exchange degeneracy”

Degenerate perturbation theory

- We must use superposition states
- $(H_0 + H') \psi = (E_0 + \Delta E) \psi$
 - where:
 - $\psi = \alpha \psi_{1s}(\vec{r}_1) \psi_{nl}(\vec{r}_2) + \beta \psi_{nl}(\vec{r}_1) \psi_{1s}(\vec{r}_2)$
 - $E_0 = E_{1s} + E_{nl}$

$$H' \psi = \Delta E \psi$$

$$H' \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = \Delta E \begin{pmatrix} \alpha \\ \beta \end{pmatrix}$$

$$H' = \begin{pmatrix} J & K \\ K & J \end{pmatrix}$$

$$J = \int |\psi_{1s}(\vec{r}_1)|^2 \frac{1}{r_{12}} |\psi_{nl}(\vec{r}_2)|^2 d\vec{r}_1 d\vec{r}_2$$

$$K = \int \psi_{1s}^*(\vec{r}_1) \psi_{nl}^*(\vec{r}_2) \frac{1}{r_{12}} \psi_{1s}(\vec{r}_2) \psi_{nl}(\vec{r}_1) d\vec{r}_1 d\vec{r}_2$$

- J : the “direct integral”
 - Coulomb interaction between the two charge clouds
 - Increases energy
- K : the “exchange integral”
 - a quantum interference effect

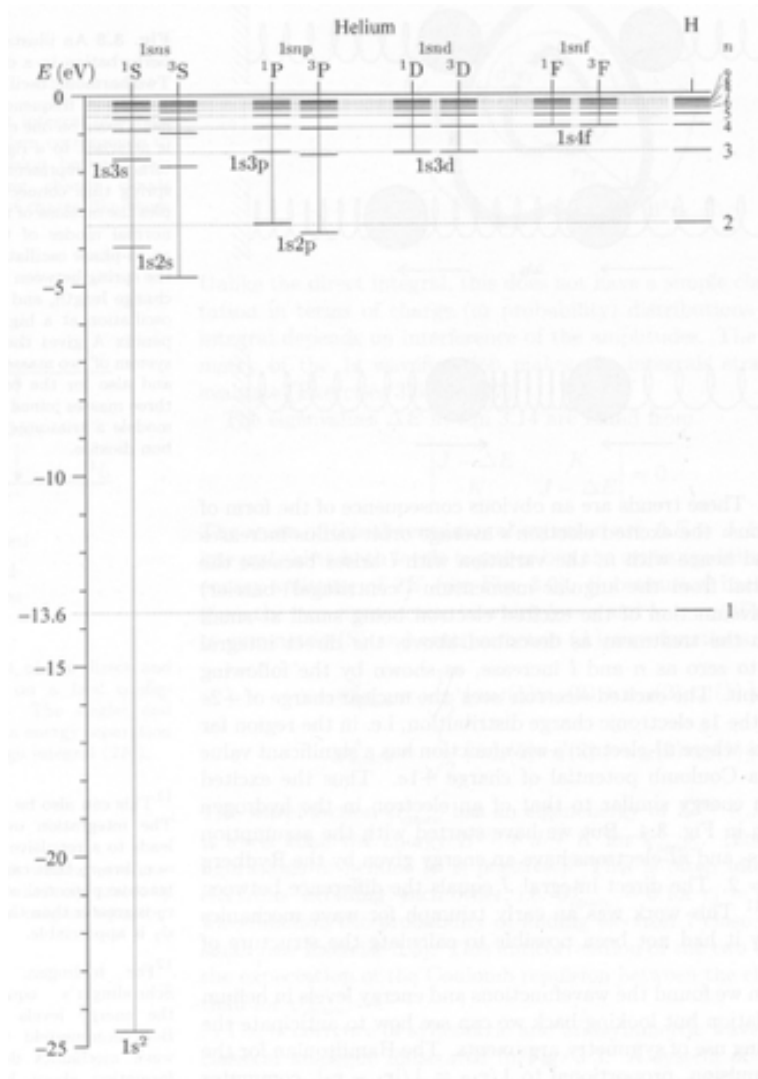
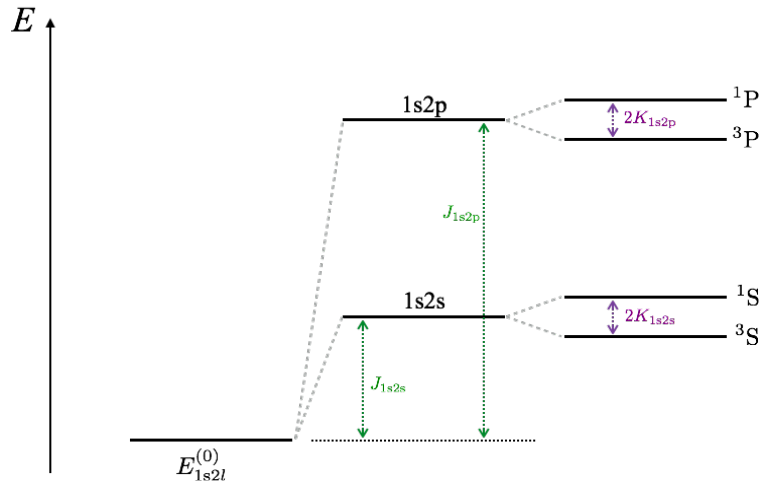
$$\Delta E = J \pm K$$

$$E_{1s nl}^{\pm} = E_{1s nl}^0 + J \pm K$$

- The wave functions are symmetric or anti-symmetric :

$$\begin{cases} \psi_+^{(0)}(\vec{r}_1, \vec{r}_2) \equiv \frac{1}{\sqrt{2}} [\psi_{1s}(\vec{r}_1) \psi_{nl}(\vec{r}_2) + \psi_{nl}(\vec{r}_1) \psi_{1s}(\vec{r}_2)] \\ \psi_-^{(0)}(\vec{r}_1, \vec{r}_2) \equiv \frac{1}{\sqrt{2}} [\psi_{1s}(\vec{r}_1) \psi_{nl}(\vec{r}_2) - \psi_{nl}(\vec{r}_1) \psi_{1s}(\vec{r}_2)] \end{cases}$$

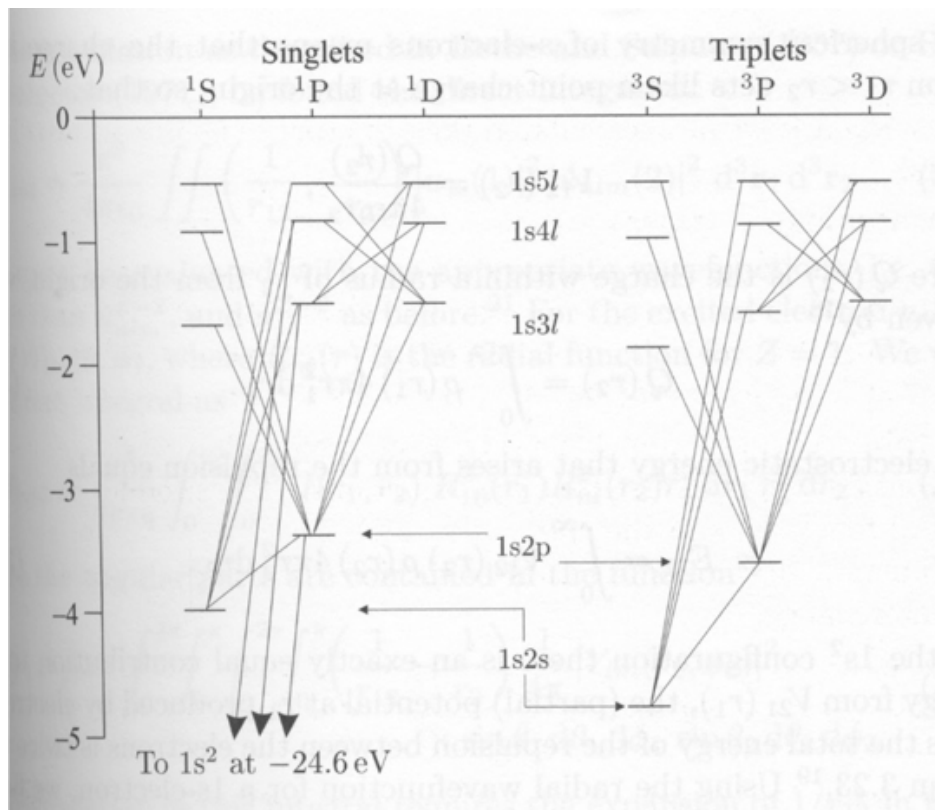
- entangled states



Transitions in He

- Selection rule for the total spin :

$$\Delta S = 0$$
- A two-electron atom will only have singlets ($S = 0$) and triplets ($S = 1$)
- There will never be transitions between a singlet and triplet



- He gives an appearance of having two separate spectra